# First demonstration of field-free perpendicular SOTMRAM for ultrafast and high-density embedded memories 

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#### Abstract

For the first time, we experimentally demonstrate the field-free switching in multi-pillar (MP) spin-orbit torque magnetic random-access memory (SOT-MRAM) devices, which are CMOS-compatible 300 mm integrated perpendicular MTJs (p-MTJs). The field-free switching (FFS) is achieved by integrating an additional in-plane (IP) magnet layer below the conventional heavy metal layer, forming a hybrid spin source layer. The in-plane magnet contributes to additional unconventional spin-orbit torque, breaking the symmetry for field-free switching and leading to high SOT switching efficiency. We demonstrate ultrafast field-free switching with current pulses down to 0.3 ns , corresponding to a power consumption of $60 \mathrm{fJ} / \mathrm{bit}$. Moreover, this FFS scheme is fully compatible with the standard integration process and the voltage-gated SOT (VG-SOT) switching in MP devices. Selective operations of independent write and read between multiple MTJs on a shared SOT track can also be achieved without external magnetic field. The FFS concept is scalable, agnostic to SOT material, and enables the reduction of the external periphery (i.e.: transistors). Thus, our proposed concept is advantageous for further improving the density and energy efficiency of SOT-MRAM technology.


## I. Introduction

MRAM has become a strong candidate for embedded memories, due to its non-volatility, infinite endurance ( $>10^{12}$ ), fast speed (ns timescale), low power consumption and good CMOS and BEOL compatibility. To date, spin-transfer torque (STT)-MRAM has become a mature technology in recent years with major foundries announcing large volume production of embedded STT-MRAM products, which have been used as the replacement of e-FLASH [1]. However, the nanoseconds incubation delay and endurance issues of STT-MRAM makes it unsuitable for high-level memory, such as L1/L2 cache SRAM replacement [2].

Recently, SOT-MRAM has been highlighted as a potential candidate owing to its sub-ns switching speeds and enhanced endurance due to separated read/write paths [2, 3]. However, the relatively high switching current and the 3-terminal configuration adversely impacts its reliability and bit-cell area, thus limiting the integration density. Moreover, an in-plane external field is typically required to achieve deterministic switching in perpendicular SOT-MRAM, which is also an obstacle for integration. To mitigate these issues, concepts of selective operations through voltage-control multi-bits and SOT shared-write-channel $[4,5,6]$ have been proposed. However, field-free solutions compatible with multi-pillar
operations are still not well developed. Although several methods have been proposed to remove the necessity of the external field $[7,8,9]$, these solutions compromise scalability and efficiency due to the complexity for the write operation and device integration, making the solutions impractical for MRAM products.

In this paper, we propose and demonstrate a field-free SOT switching scheme in multi-bit integrated p -MTJs with bitselective $\mathrm{R} / \mathrm{W}$ operations on 300 mm wafers. Our demonstration removes the necessity of the external field and reduces the number of transistors/bit and switching current simultaneously, providing a reliable solution for high-density and low power SOT-MRAM.

## II. Hybrid SOT for field-free SOT-MRAM

## A. Concept of field-free switching multi-pillar device

In conventional SOT devices, the spin currents are generated from heavy metals (such as W, Pt, and Ta) or their alloys, exhibiting an in-plane spin polarization [10]. To switch a p-MTJ device, an in-plane external field is generally needed to break the symmetry for achieving deterministic switching. Magnets with in-plane anisotropy can also contribute to unconventional SOT with out-of-plane spin polarization for FFS [11]. Here, we propose an FFS scheme of a hybrid SOT track, which consists of a heavy metal and an in-plane magnet, combining the high spin generation and unconventional SOT for FFS (as shown in Figure 1). In this scheme, the additional in-plane magnet layer is separated from the MTJ pillars by the SOT track, which can accommodate multiple pillars, thereby enabling voltage-controlled selective $\mathrm{R} / \mathrm{W}$ operations.

## B. Film stack development and device integration

In this work, each MTJ consists of a top-pinned MTJ stack defined as substrate/[magnet/W]/CoFeB (FL)/MgO/CoFeB (RL). Compared to the conventional SOT-MTJ structure, an additional magnetic layer was deposited below the W layer, forming a hybrid SOT track to generate the spin current and break the symmetry. Figure 2 shows the resistivity of W in a magnet/W $(t \mathrm{~nm})$ structure as deposited and after annealing. The high resistivity of $\sim 130 \mu \Omega \cdot \mathrm{~cm}$ is obtained for W thickness less than 3.2 nm , indicating a $\beta-\mathrm{W}$ [3].

The devices were fabricated in the IMEC 300 mm pilot line. Figure 3 shows the cross sections views of the proposed MTJ devices. The in-plane magnet below the heavy metal has minimal impact on the integration process, which is the same as in our previous reports [3,5]. In Fig. 4, we show a typical hysteresis loop of a 60 nm -sized MTJ. The square shape of the
$R-B$ curve indicates good PMA of the device. The size dependent MTJ performance was evaluated (over 100 devices/CD size) as shown in Fig. 5. For 80 nm MTJs, the median values of TMR, coercivity ( $B_{\mathrm{c}}$ ) and offset field ( $B_{\text {offset }}$ ) are $105 \%$, 55.4 mT , and -22.3 mT , respectively. In Fig. 6, the thermal stability $(\Delta)$ and anisotropy field $\left(B_{k}\right)$ of FFS samples are $\sim 50 k_{\mathrm{b}} \mathrm{T}$ and 120 mT , respectively. No distinct differences in $\Delta$ and $B_{\mathrm{k}}$ were found between the FFS and W-POR samples. The SOT track resistivity is $155 \mu \Omega \cdot \mathrm{~cm}$, which is in line with the resistivity of $\beta-\mathrm{W}$. The bottom magnetic layer has limited impact on the free and SOT layer.

## III. FIELD-FREE SOT SWTICHING PEFORMANCE

## A. Field-free SOT switching

Figure 7 shows the typical current-induced magnetization switching curves in an 80 nm MTJ device under different inplane external fields $\left(B_{x}\right)$. The deterministic switching loops can be obtained even at zero $B_{x}$. The field-free switching is achieved in devices with different CDs. We extracted the effective fields ( $B_{\text {eff }}$ ) for FFS from the sign change in switching polarity. The results are summarized in Fig. 8. $B_{\text {eff }}$ increases with shrinking the CD size, which is beneficial for device scalability. In small devices (with CD $<100 \mathrm{~nm}$ ), strong $B_{\text {eff }}$ is obtained, which can go up to 40 mT . While in large MTJ devices, it is difficult to get full field-free switching due to insufficient $B_{\text {eff. }}$

## B. SOT efficiency and switching performance

The SOT efficiency is characterized by current-induced hysteresis loop shift measurement [12]. From Fig. 9, the nonzero SOT efficiency $\left(\chi_{\text {SOT }}\right)$ at $B_{x}=0 \mathrm{mT}$ confirms the unconventional SOT for field-free switching, corresponding to an effective spin Hall angle $\left(\theta_{\mathrm{SH}}\right)$ of $-18.1 \%$. A total effective $\theta_{\text {SH }}$ of $-37.6 \%$ in hybrid SOT track can be obtained under $B_{x}=$ -30 mT (green dash line $\theta_{\mathrm{SH}, \beta-\mathrm{w}}=-32 \%$ [3]), indicating that the in-plane magnet contributes the additional spin generation (with $\theta_{\mathrm{SH}}=-5.6 \%$ ), which is in good agreement with previously reported results [13].

Figure 10 shows the FFS loops for different pulse widths. The ultrafast switching can be obtained with pulse down to 0.3 ns . In a typical 60 nm device, the switching current density at 1 ns is $5.7 \times 10^{11} \mathrm{~A} / \mathrm{m}^{2}$ with a power consumption of 60 fJ . The benchmark plots of $J_{\text {sw }}$ vs. $\Delta$ and $E_{\text {sw }}$ vs. $\Delta$ are shown in Figure 11. Devices with hybrid SOT track show high SOT efficiency, which is comparable with W-POR samples.

Figure 12 shows the $63 \%$-lifetime measurement results of FFS devices. From extrapolation [14], an endurance of $>10^{12}$ could be expected in both $\mathrm{CD}=60$ and 80 nm devices at the switching current for $P_{\mathrm{sw}}=50 \%$.

## IV. VolTAGE-GATED FIELD-FREE SWITCHING

## A. VCMA effect and $V G$-SOT switching in single pillar

To further enhance the SOT efficiency and device density, a voltage-gated scheme was proposed in the FFS devices by utilizing the VCMA effect $[4,5]$. We evaluate the impact of gate voltage $\left(V_{\mathrm{G}}\right)$ for VCMA-induced modulations of magnetic anisotropy $\left(B_{k}\right)$ and switching current $\left(I_{\mathrm{sw}}\right)$ in our FFS devices (as illustrated in Figure 13). Figure 14 shows the $R-B$ loops with
$V_{\mathrm{G}}$. We extracted the anisotropy field $B_{k}$ from hysteresis loops of MTJs with $V_{\mathrm{G}}$. $B_{k}-V_{\mathrm{G}}$ curves of MTJs show a slope of -34.7 $\mathrm{mT} / \mathrm{V}$ in Fig. 15. The corresponding VCMA coefficient is estimated to be $\sim 17.5 \mathrm{fJ} / \mathrm{Vm}$.

Figure 16 shows the switching loop and switching probability $\left(P_{\mathrm{sw}}\right)$ for different $V_{\mathrm{G}}$ in a 60 nm MTJ at 1 ns pulse. $I_{\mathrm{sw}}$ shows a linear dependence on $V_{\mathrm{G}}$. A $\sim 30 \%$ reduction in $I_{\mathrm{sw}}$ can be obtained by applying $V_{\mathrm{G}}$ within 0.5 V .

## B. Voltage-gated FFS and selective operation in multi-pillar

 devicesSimilarly, we performed the VG-SOT switching measurements in the MP devices. In the measurements, we completely removed the magnet fields, the deterministic switching with different gate voltages can also be obtained. Figure 17 shows the typical VG-SOT switching loops in one of four MTJs in MP devices. By applying gate voltages, the FFS could be modulated to a higher/lower critical switching voltage, enabling the selective operations of multi-bits [5].

Combined on our experimental observations, the results on FFS SOT performance with previously reported technologies are summarized in Table 1. The proposed FFS scheme by hybrid SOT track exhibits good compatibility with high-density integration and maintains high SOT efficiency. Further optimizations may focus on improving switching efficiency, reliability (WER performance) and VCMA coefficient.

## V. CONCLUSIONS

We demonstrated a field-free perpendicular SOT-MRAM design by a hybrid spin source layer, consisting of a heavy metal and an in-plane magnet. The hybrid spin source enables ultrafast field-free switching with high switching efficiency and fast operational speed down to 0.3 ns . Moreover, this scheme is suitable for multi-bit SOT devices and is CMOS and back-end-of-line compatible. Our solution provides a path for high performance and high-density MRAM applications.

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## REFERENCES

[1] J. G. Alzate et al., IEEE IEDM 2.1.1-2.4.4 (2019). [2] Cubukcu, M. et al., IEEE Trans. Magn. 54, 9300204 (2018). [3] K. Garello et al., IEEE Symp. On VLSI Circuit, 81-82 (2018). [4] H. Yoda et al., IEEE IEDM 27.6.1-27.6.4 (2016). [5] K. Cai et al., IEEE Symp. On VLSI Technology, 375-376 (2022). [6] S.Z. Peng et al., IEEE IEDM 28.6.1-28.6.4 (2019).[7] M. Wang et al., Nature Electronics, 1, 582-588 (2018). [8] D.Q. Zhu et al. IEEE IEDM 17.5.117.5.4 (2021). [9] K. Garello et al., IEEE Symp. On VLSI Technology, 194-195 (2019). [10] I. M. Miron et al., Nature 476, 189-193 (2011). [11] S.H.C. Baek et al, Nature Materials 17, 509-513 (2018). [12] C. Pai et al., Phys. Rev. B 93,144409 (2016). [13] W.L. Yang et al., Appl. Phys. Lett. 120,122402 (2022). [14] S. Van Beek et al., IEEE IRPS, 4A.2-1-4A.2-4 (2022). [15] H. Honjo et al., IEEE IEDM 28.5.1-28.5.4 (2019). [16] M.Y. Song et al., IEEE Symp. On VLSI Technology, 377-378 (2022).


Fig. 1. Schematic of the field-free switching MTJ stack with hybrid SOT track design: SOT(magnet + heavy metal)/ $\mathrm{CoFeB}(\mathrm{FL}) / \mathrm{MgO} / \mathrm{CoFeB}(\mathrm{RL}) / \mathrm{SAF}(\mathrm{HL})$.


Fig. 3. (a) An optical image of MTJ devices. (b) TEM cross section view of a single pillar device. (c) Zoom in view of a 60 nm MTJ. (d) the integrated multi-pillar SOT device.


Fig. 4. Hysteresis loops of 60 nm FFS-SOT and W-POR (reference stack) devices. Bottom in-plane magnet shows low impact on FL and RL.


Fig. 5. Statistics of basic characterizations on FFS-SOT MTJ samples (over 100 devices/CD size). (a) TMR, (b) $\mathrm{B}_{\mathrm{c}}$, (c) $B_{\text {offset }}$


Fig. 6. Comparisons between FFS-SOT and W-POR samples. (a) thermal stability ( $\Delta$ ); (b) anisotropy field $\left(B_{\mathrm{k}}\right)$.


Fig. 7. Current-induced switching loops of an 80 nm MTJ device for different $B_{x}$. As indicated, the full switching loop is obtained for $B_{\mathrm{x}}=0 \mathrm{mT}$. The switching diminishes at 35 mT when $B_{x}$ fully compensates the effective field.


Fig. 8. Extracted effective field $B_{\text {eff }}$ for field-free switching of MTJs with different CD sizes.


Fig. 9: Damping-like spin-orbit torque efficiency of hybrid SOT track at $B_{x}=0$ and -30 mT . The green dash line shows the reference results of W-POR.
(c)

Fig. 10. Ultrafast field-free SOT switching. (a) Switching loops for different pulse widths without external magnetic field. (b) Critical switching current as a function of the inverse of pulse width. (c) switching energy as a function of pulse width. It shows the sub-ns switching and low energy consumption.


Fig. 11. (a) Switching current density and (b) energy consumption at 1 ns as a function of retention ( $\Delta$ ) in W-POR and FFS devices. The red points show the results of 60 nm and 80 nm devices. The hybrid SOT track shows comparable efficiency as W-POR samples.


Fig. 13. Schematic of voltage controlled magnetic anisotropy (VCMA) effect and multi-pillar devices. (a) the gate voltage modulates the energy barrier for switching. (b) the hybrid SOT track enables the field-free operations on multi-pillar devices.

Fig. 14: Hysteresis loops of a 60 nm FFS-MTJ devices for different gate voltages.


Fig. 15. Anisotropy field $B_{\mathrm{k}}$ as a function of gate voltage $V_{\mathrm{g}}$ of two 80 nm FFS-MTJ devices.




Fig. 12. $63 \%$ lifetime analysis with 1 ns bipolar pulses. 14 devices stressed at each condition. Star symbols depict lifetime at critical $I_{\mathrm{sw}}$ condition (at $P_{\mathrm{sw}}=50 \%$ ).

Fig. 16. Voltage-gated SOT (VG-SOT) switching. (a) Normalized switching loops for different $\mathrm{V}_{\mathrm{G}}$ by 1 ns pulses in a 60 nm FFS device. (b) Switching probability $P_{\text {SW }}$ curves for different $\mathrm{V}_{\mathrm{G}}$. (c) Critical switching current $I_{\mathrm{sw}}$ (at $P_{\mathrm{sw}}=50 \%$ ) as a function of gate voltage $\mathrm{V}_{\mathrm{G}}$.


Fig. 17. A typical field-free VG-SOT switching by 100 ns current pulses in one of four MTJs in multipillar devices.

| Specification |  | $\begin{gathered} \text { IMEC } \\ \text { W-POR [3] } \end{gathered}$ | $\begin{gathered} \text { IMEC } \\ \text { MHM [9] } \end{gathered}$ | Tohoku Univ. Canted-SOT [15] | Beihang Univ. EB [8] | $\begin{gathered} \text { TSMC } \\ \text { Type-Y [16] } \end{gathered}$ | This work |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $$ | Type of MTJ | PMA | PMA | IMA | PMA | IMA | PMA |
|  | $\mathrm{CD}(\mathrm{nm})$ | 60 | 60 | 88*340 | 700 | 75*230 | 60 |
|  | RA $\left(\Omega \cdot \mu \mathrm{m}^{2}\right)$ | 40 | 40 | $\sim 500$ | $\sim 3000$ | 10 | 100 |
|  | TMR (\%) | 110 | 110 | 167 | 101 | 140 | 105 |
|  | $\mathrm{B}_{\mathrm{c}}(\mathrm{mT})$ | 94.5 | 83.5 | ~13.5 | N.A. | $\sim 20$ | 53.7 |
|  | $\Delta\left(\mathrm{k}_{\mathrm{b}} \mathrm{T}\right)$ | 51 | 48 | 70 | N.A. | 152 | 35 |
|  | VCMA \& MP compatible | N.A. | No | N.A. | N.A. | N.A. | Yes |
| $\begin{aligned} & \text { Wi } \\ & 0 . \\ & 0 \\ & 0 \\ & \hline \end{aligned}$ | Annealing ( ${ }^{\circ} \mathrm{C}$ ) | 300 | 300 | 400 | 300 | 400 | 300 |
|  | Wafer Process | 300 mm | 300 mm | 300 mm | N.A. | N.A. | 300 mm |
|  | External Field (mT) | 32 | 0 | 0 | 0 | 0 | 0 |
|  | $J_{\mathrm{c}}\left(\mathrm{MA} / \mathrm{cm}^{2}\right) @ 1 \mathrm{~ns}$ | 120 | 126 | 23.6 | 170 @10ns | 61 | 57 |
|  | Pulse width (ns) | 0.21 | 0.30 | 0.35 | 10 | 0.4 | 0.30 |
|  | Endurance | $>5 \times 10^{11}$ | $>10^{11}$ | N.A. | $>10^{10}$ | $>7 \times 10^{12}$ | $>10^{12}$ |

Table 1: Benchmark table of our FFS devices with hybrid SOT track and other published field-free SOT devices.

